

Large Amplitude Motion Effects in the TPE Spectrum of Methyl Isocyanate

O. J. Harper,^a B. Gans,^a S. Boyé-Péronne,^a L. H. Coudert,^a
J.-C. Loison,^b G. A. Garcia,^c and J.-C. Guillemin^d

^aISMO, Universités Paris-Sud & Paris-Saclay, Orsay, France

^bInstitut des Sciences Moléculaires, Université de Bordeaux, Talence, France

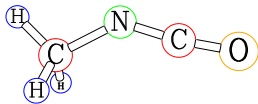
^cDESIRS beamline, Synchrotron SOLEIL, Gif-sur-Yvette, France

^dENSC, Univ. Rennes, Rennes, France

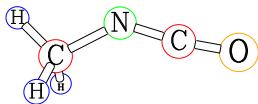
ISMS, June 17–21, 2019



Outline

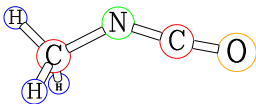


Outline



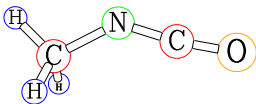
- Microwave & photoelectron spectra

Outline



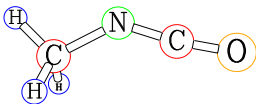
- Microwave & photoelectron spectra
- Threshold photoelectron spectrum (TPES) recorded in this work

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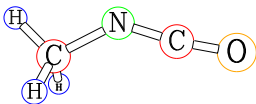
- Microwave & photoelectron spectra
- Threshold photoelectron spectrum (TPES) recorded in this work
- How should this spectrum be modeled

Outline



- Microwave & photoelectron spectra
- Threshold photoelectron spectrum (TPES) recorded in this work
- How should this spectrum be modeled
- Comparisons between observed and calculated TPES

The microwave spectrum

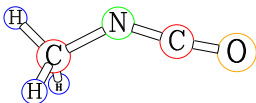


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²Koput, *J. Molec. Spectrosc.* **106** (1984) 12

³Cernicharo, Kisiel, Tercero *et al.*, *A&A* **587** (2016) L4

The microwave spectrum



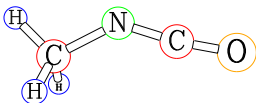
- MW spectrum first recorded in the sixties.^{1,2} New measurements carried out recently.³ A dense spectrum was observed

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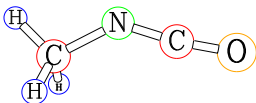
- MW spectrum first recorded in the sixties.^{1,2} New measurements carried out recently.³ A dense spectrum was observed
- Two large amplitude motions need to be accounted for to reproduce the frequency of *a*-type lines²

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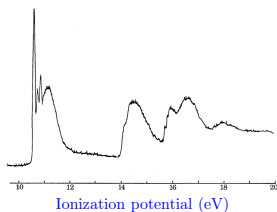
- MW spectrum first recorded in the sixties.^{1,2} New measurements carried out recently.³ A dense spectrum was observed
- Two large amplitude motions need to be accounted for to reproduce the frequency of *a*-type lines²
- Detected in the Orion clouds³

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The photoelectron spectrum

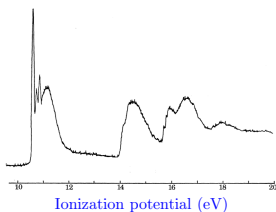


¹Eland, *Phil. Trans. Roy. Soc. Lond. A* **268** (1970) 87

²Cradock, Ebsworth, & Murdoch, *J. Chem. Soc. Faraday Trans. II* **68** (1972) 86

³Pasinszki, Vezsprèmi, Fehèr, Kovač, Klasinc, & McGlynn, *Int. J. Qant. Chem.* **26** (1992) 443

The photoelectron spectrum



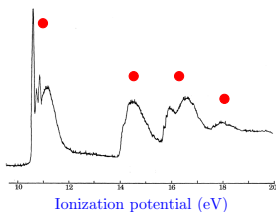
- The **photoelectron** spectrum has been investigated and spectra spanning the energy range from **10** to **18 eV** were recorded¹⁻³

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The photoelectron spectrum



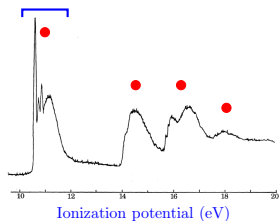
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- **Four** electronic bands were observed

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The photoelectron spectrum



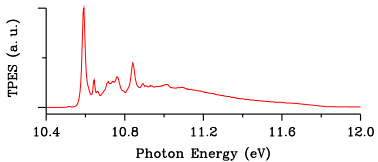
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- **Four** electronic bands were observed
- In this talk, the low energy band, from **10.4** to **12 eV**, is investigated with a higher resolution

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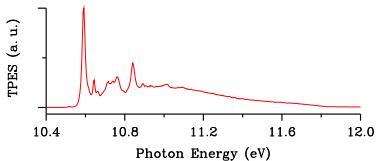
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The TPES recorded in this work

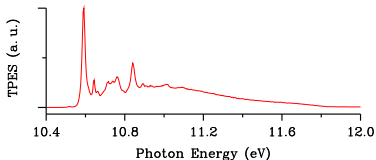


The TPES recorded in this work



- A TPES was recorded using [synchrotron](#) radiation

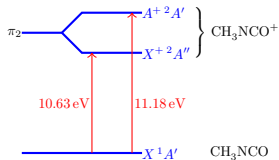
The TPES recorded in this work



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- Energy [difference](#) between cation & neutral levels

$$h\nu = E(\text{CH}_3\text{NCO}^+) - E(\text{CH}_3\text{NCO})$$

Spectrum modeling

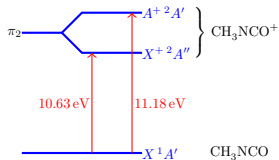


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Spectrum modeling



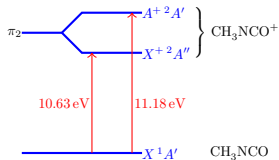
- The $\left\{ \begin{array}{l} X^+2A'' \leftarrow X^1A' \\ A^+2A' \leftarrow X^1A' \end{array} \right.$ ionizing transitions were observed^{1,2}

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Spectrum modeling



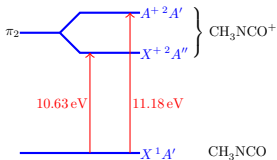
- The $\left\{ \begin{array}{l} X^+2A'' \leftarrow X^1A' \\ A^+2A' \leftarrow X^1A' \end{array} \right.$ ionizing transitions were observed^{1,2}
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Spectrum modeling



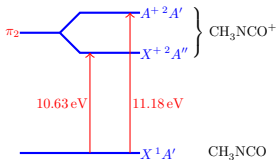
- The $\begin{cases} X^+2A'' \leftarrow X^1A' \\ A^+2A' \leftarrow X^1A' \end{cases}$ ionizing transitions were observed^{1,2}
- Energy levels of the neutral and the cation
- Eigenvectors also needed to compute line intensities³

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Spectrum modeling



- The $\begin{cases} X^+2A'' \leftarrow X^1A' \\ A^+2A' \leftarrow X^1A' \end{cases}$ ionizing transitions were observed^{1,2}
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- Eigenvectors also needed to compute line intensities³
- The π_2 orbital is located at the NCO group^{1,2}

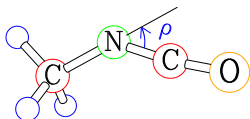
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The 2 large amplitude motions

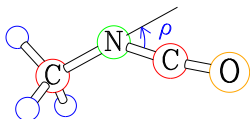
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Bending mode

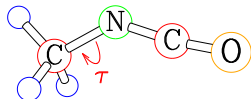
$$t = \cos(\pi - \rho)$$

The 2 large amplitude motions



Bending mode

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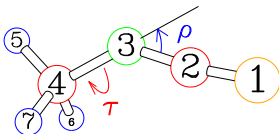


Torsion

$$\tau$$

Energy level calculation: Kinetic energy

A reference configuration $\mathbf{a}(\rho, \tau)_i$, with $1 \leq i \leq 7$, is chosen

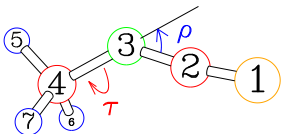


¹Kręglewski & Jensen, *J. Molec. Spectrosc.* **103** (1984) 312

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Bending-Torsion-Rotation energy levels are calculated with a 5-dimensional Hamiltonian^{1,2}

$$H(t, \tau, \Omega)$$

using the 5×5 generalized inverse inertia tensor $\boldsymbol{\mu}(t, \tau)$

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Energy level calculation: Potential energy surface

The bidimensional potential energy surface is written

$$V(\rho, \tau) = V_0(\rho) + \frac{1}{2}\rho V_3^1 \cos 3\tau$$

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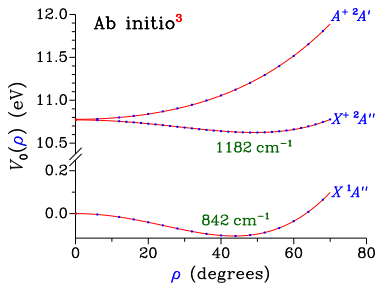
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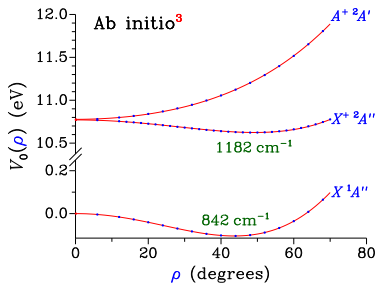
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State	V_3^1/cm^{-1}
X^1A'	20, ¹ 16.2, ² 44.5 ³
X^+2A''	235.5 ³
A^+2A'	?

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Energy level calculation: Basis set functions

Bending-torsion-rotation basis set functions

$$(1-t)^{\alpha/2}(1+t)^{\beta/2}P_n^{(\alpha,\beta)}(t) \exp(im\tau)/\sqrt{2\pi} \quad |J, k, M\rangle$$

¹Coudert, *J. Molec. Spectrosc.* **104** (1989) 123

Energy level calculation: Basis set functions

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Singularity at the linear $t = -1$ configuration

$$\mu(t)_{zz}, \mu(t)_{\tau z}, \text{ and } \mu(t)_{\tau\tau} \simeq \frac{1}{1-t^2}$$

accounted for taking appropriate values¹ for α and β

¹Coudert, *J. Molec. Spectrosc.* **104** (1989) 123

Energy level calculation: Results

Calculated $J = 1$ bending-torsion-rotation energies for the neutral

A_1					A_2					E				
v_b	k	m	Cal ¹	Cal ²	v_b	k	m	Cal ¹	Cal ²	v_b	k	m	Cal ¹	Cal ²
0	1	0	4.15	4.15	0	0	0	0.00	0.00	0	-1	1	5.13	5.13
0	-1	3	61.17	61.17	0	1	0	4.15	4.15	0	0	1	8.71	8.70
0	0	3	80.17	80.17	0	-1	3	61.17	61.17	0	1	1	18.95	18.95
0	1	3	107.16	107.16	0	0	3	80.46	80.46	0	-1	2	26.51	26.51
1	1	0	196.22	196.18	0	1	3	107.16	107.16	0	0	2	36.33	36.33
					1	0	0	191.34	191.30	0	1	2	55.75	55.75
					1	1	0	196.22	196.18	0	-1	4	115.34	115.29
										0	0	4	141.77	141.73
										0	1	4	175.46	175.41
										0	-1	5	186.71	186.67
										1	-1	1	196.49	196.45
										1	0	1	200.81	200.77

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²This work

The $X^+ 2A'' \leftarrow X^1A'$ ionizing transition @ 10K

Selection rules for rotational components $\left\{ \begin{array}{l} |\Delta J| \leq 2, |\Delta k| = 1 \\ A_{1,2} \leftrightarrow A_{1,2} \text{ or } E \leftrightarrow E \end{array} \right.$

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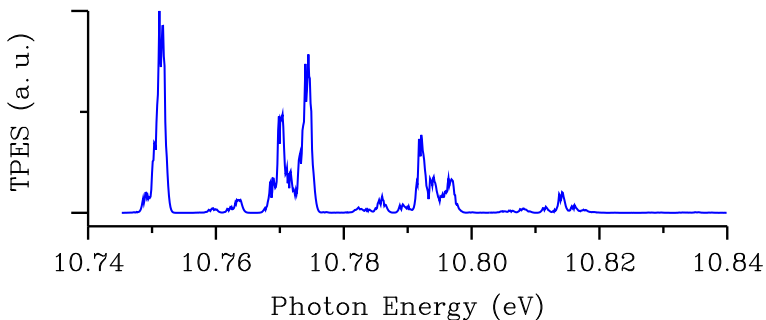
e^- are removed¹ from the π_2 orbital with $\left\{ \begin{array}{l} l = 2 \\ \lambda = 1 \end{array} \right.$

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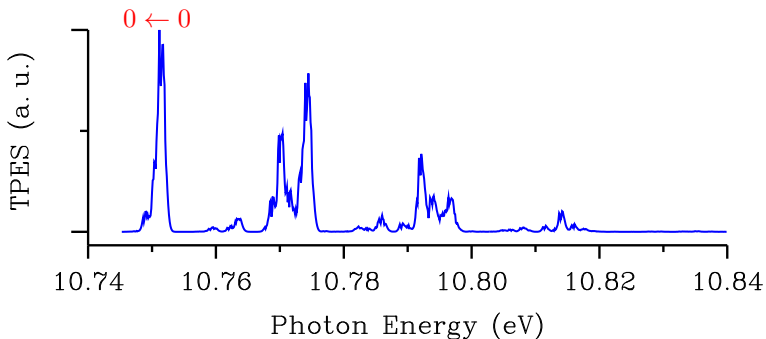


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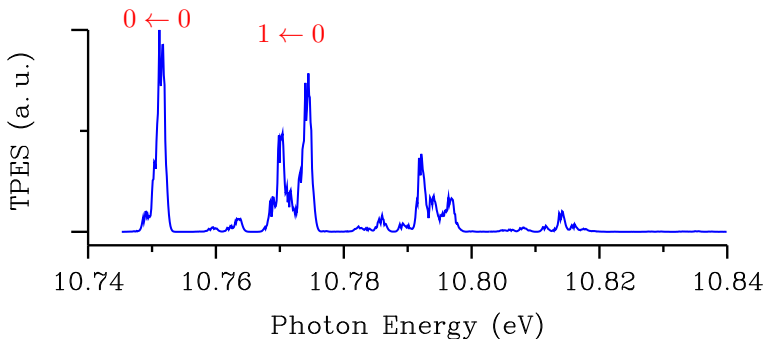


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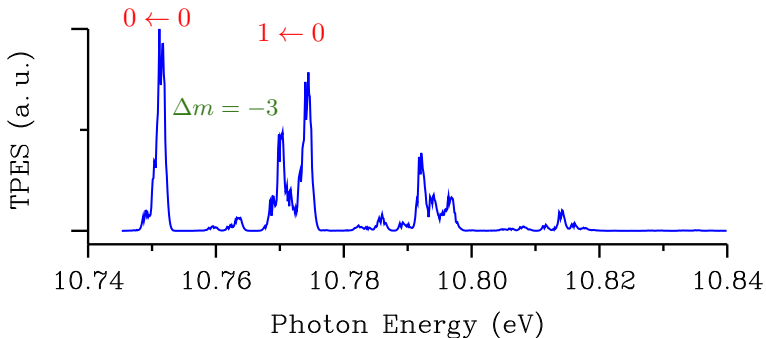


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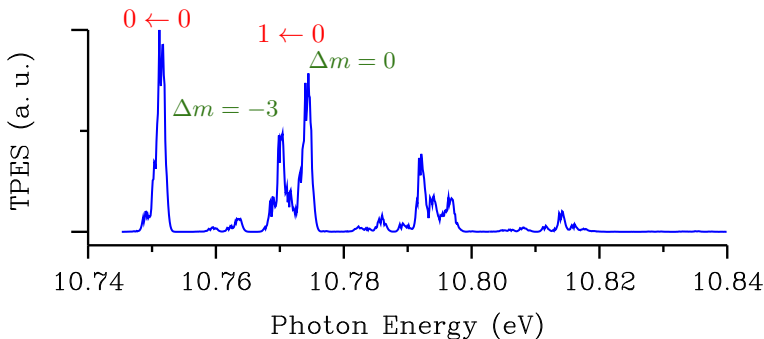


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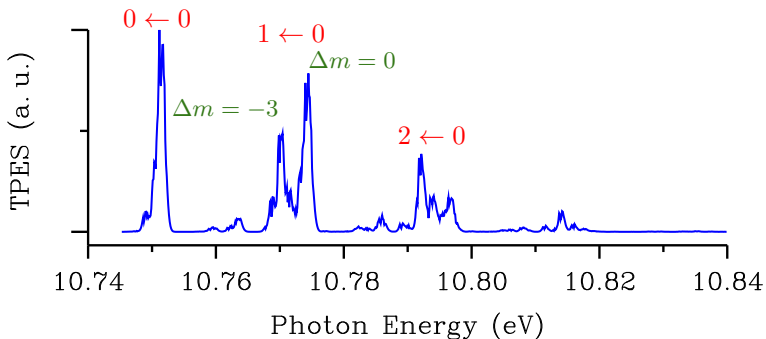


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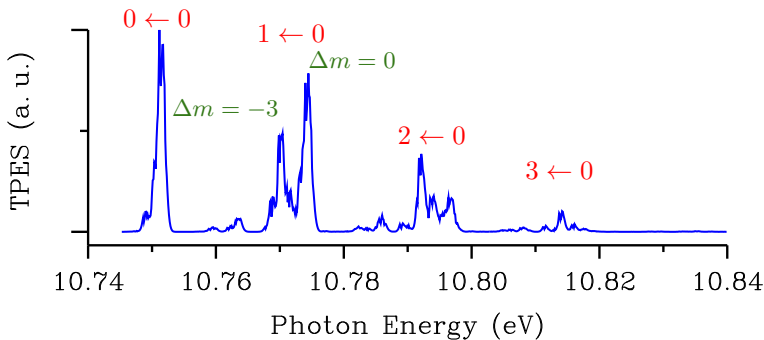


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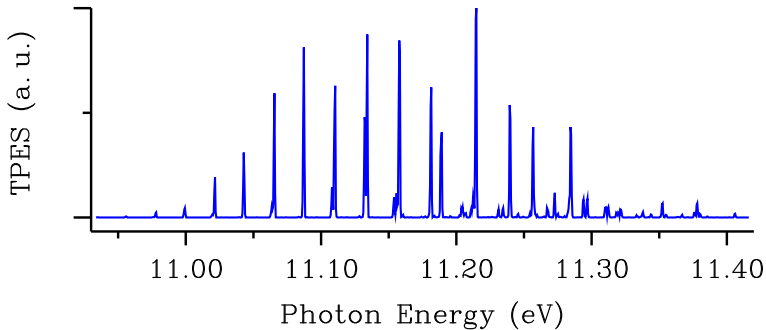
Selection rules for rotational components $\left\{ \begin{array}{l} |\Delta J| \leq 2, |\Delta k| = 1 \\ A_{1,2} \leftrightarrow A_{1,2} \text{ or } E \leftrightarrow E \end{array} \right.$

e^- are removed¹ from the π_2 orbital with $\left\{ \begin{array}{l} l = 2 \\ \lambda = 1 \end{array} \right.$



¹Willitsch & Merkt, *Int. J. Mass Spectrosc.* **245** (2005) 14

The $A^+ 2A'' \leftarrow X^1A'$ ionizing transition @ 10K



TPES calculation

Franck Condon progressions due to **small amplitude** vibrational modes involving the **NCO** group^{1,2} should be taken into account

¹Hirschmann, Kniseley, & Fassel, *Spectrochim. Acta* **21** (1965) 2125

²Pasinszki, Vezsprèmi, Fehèr, Kovač, Klasinc, & McGlynn, *Int. J. Qant. Chem.* **26** (1992) 443

TPES calculation

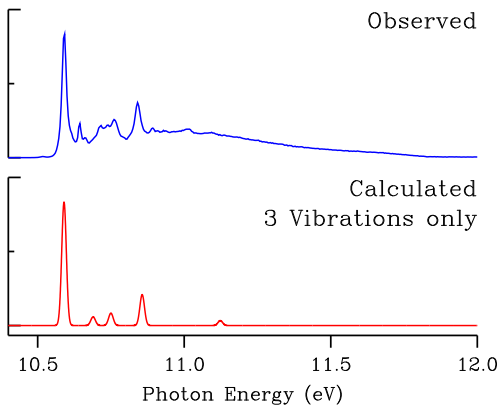
Franck Condon progressions due to **small amplitude** vibrational modes involving the **NCO** group^{1,2} should be taken into account

Mode	X^1A'	X^+2A''
C-N	852	800
$\nu_s(\text{NCO})$	1437	1290
$\nu_a(\text{NCO})$	2288	2150

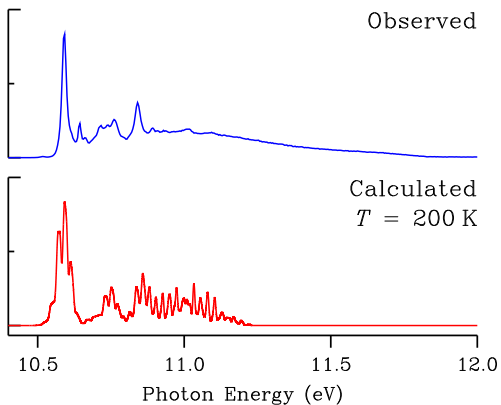
¹Hirschmann, Kniseley, & Fassel, *Spectrochim. Acta* **21** (1965) 2125

²Pasinszki, Vezsprèmi, Fehèr, Kovač, Klasinc, & McGlynn, *Int. J. Qant. Chem.* **26** (1992) 443

Observed and calculated TPES (3 vibrations only)



Observed and calculated TPES

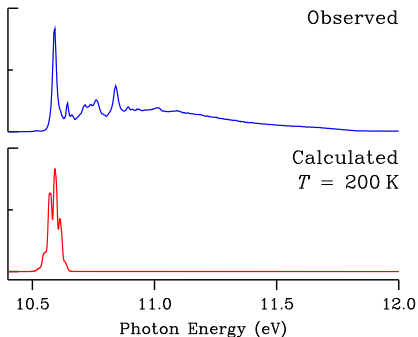


Observed and calculated TPES

Without the Franck Condon progressions due to small amplitude vibrational modes and the $A^+ 2A' \leftarrow X 1A'$ ionizing transition

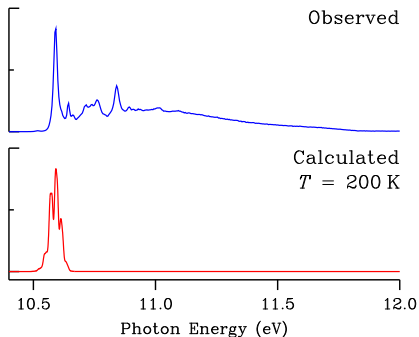
Observed and calculated TPES

Without the Franck Condon progressions due to small amplitude vibrational modes and the $A^+ 2A' \leftarrow X^1A'$ ionizing transition



Observed and calculated TPES

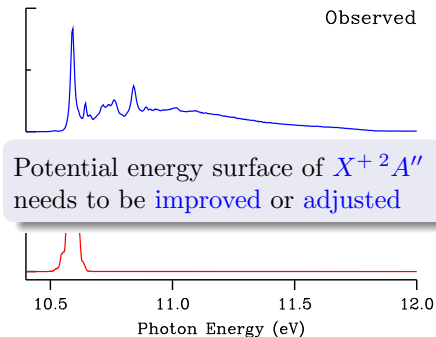
Without the Franck Condon progressions due to small amplitude vibrational modes and the $A^+ 2A' \leftarrow X^1A'$ ionizing transition



The $X^+ 2A'' \leftarrow X^1A'$ ionizing transition leads to a feature which is either too **wide** or too **narrow**

Observed and calculated TPES

Without the Franck Condon progressions due to small amplitude vibrational modes and the $A^{+2}A' \leftarrow X^1A'$ ionizing transition



The $X^{+2}A'' \leftarrow X^1A'$ ionizing transition leads to a feature which is either too wide or too narrow

Thank You