

On the  
usefulness of  
electron  
propagator  
methods for a  
reliable  
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experimental  
observables

Lorenzo Paoloni

Electron  
propagator  
methods

Computational  
simulation of  
vibrationally  
resolved  
ultraviolet  
photoelectron  
spectra

# On the usefulness of electron propagator methods for a reliable computation of experimental observables

Lorenzo Paoloni<sup>1</sup>

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(Italy)



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# Overview of this presentation

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- ▶ Electron propagator methods
- ▶ Computational simulation of vibrationally resolved ultraviolet photoelectron spectra (UPS)
  - Description of the computational protocol
  - Application to 6 organic molecules\*
- ▶ From gas phase to solutions: **missing part**
- ▶ Acknowledgements

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\*Published in:

L. Paoloni, M. Fusé, A. Baiardi, V. Barone, *J. Chem. Theory Comput.*, **2020**, 16, 5218-5226.

# Electron propagator theory: the spectral representation

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Electron propagator methods

Computational simulation of vibrationally resolved ultraviolet photoelectron spectra

- ▶ A unified treatment of ionization processes
  - Electron detachment:  $A \longrightarrow A^+ + e^-$
  - Electron attachment:  $A + e^- \longrightarrow A^-$
- ▶ A feasible computational approach to systematically improve estimates based on Koopmans' theorem

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See:

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Propagators (or double time green's functions) are often studied in terms of their spectral representation

$$G_{pq}(E) = \langle \langle a_p^\dagger; a_q \rangle \rangle_E = \lim_{\eta \rightarrow +0} \left[ \sum_n \frac{\langle \phi_0^{N_{el}} | a_p^\dagger | \phi_n^{N_{el}-1} \rangle \langle \phi_n^{N_{el}-1} | a_q | \phi_0^{N_{el}} \rangle}{E + E_n^{N_{el}-1} - E_0^{N_{el}} - i\eta} + \sum_m \frac{\langle \phi_0^{N_{el}} | a_q | \phi_m^{N_{el}+1} \rangle \langle \phi_m^{N_{el}+1} | a_p^\dagger | \phi_0^{N_{el}} \rangle}{E - E_m^{N_{el}+1} + E_0^{N_{el}} + i\eta} \right].$$

- ▶  $G_{pq}(E)$  is an element of the matrix  $\mathbf{G}(E)$

$p, q \equiv$  spin-orbital indices

$a_p^\dagger \equiv$  destruction operator

$a_q \equiv$  creation operator

$\phi_0^{N_{el}} \equiv$  electronic ground state of the  $N_{el}$ -electron system

$\phi_n^{N_{el}-1} \equiv$  electronic state n of the  $(N_{el} - 1)$ -electron system

$\phi_m^{N_{el}+1} \equiv$  electronic state m of the  $(N_{el} + 1)$ -electron system

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- ▶  $G_{pq}(E)$  is an element of the matrix  $\mathbf{G}(E)$

Poles correspond to  $E = E_n^{N_{el}+1} - E_0^{N_{el}}$  (electron affinities)  
and  $E = E_0^{N_{el}} - E_n^{N_{el}-1}$  (ionization potentials)

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# Electron propagator methods: the self-energy

On the usefulness of electron propagator methods for a reliable computation of experimental observables

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Electron propagator methods

Computational simulation of vibrationally resolved ultraviolet photoelectron spectra

- ▶ the inverse of the matrix  $\mathbf{G}(E)$  can be written as follows:

$$\mathbf{G}^{-1}(E) = E\mathbf{I} - \mathbf{F}^{gen} - \boldsymbol{\sigma}(E)$$

- $\mathbf{F}^{gen}$  is a generalized Fock matrix ( $\mathbf{F}^{gen} = \mathbf{F} + \boldsymbol{\Sigma}(\infty)$ )
- $\boldsymbol{\sigma}(E)$  is the (energy-dependent) self-energy term

- ▶  $\det \mathbf{G}(E)$  diverges when  $\det \mathbf{G}^{-1}(E) = 0$

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\*For example:

outer valence Green's function (OVGF), partial third-order (P3) or diagonal second-order (D2) approximations

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i.e. when the energy  $E$  is equal to a pole ( $E_{pole}$ ):

$$[\mathbf{F}^{gen} + \boldsymbol{\sigma}(E_{pole})]\mathbf{C}(E_{pole}) = E_{pole}\mathbf{C}(E_{pole})$$

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The method depends on the choice of  $\boldsymbol{\Sigma}(E) = \boldsymbol{\Sigma}(\infty) + \boldsymbol{\sigma}(E)$

- ▶  $\boldsymbol{\Sigma}(E) = 0$ : Hartree-Fock approximation
- ▶  $\Sigma_{pq}(E) = 0$  if  $p \neq q$ : diagonal approximations\*
- ▶ Otherwise: non - diagonal approximations

Diagonal approximations are often a good compromise between accuracy and computational cost

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\*For example:

outer valence Green's function (OVGF), partial third-order (P3) or diagonal second-order (D2) approximations



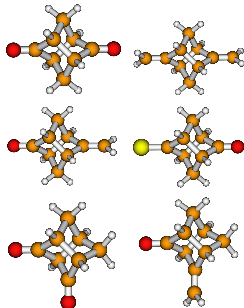
# Computational simulation of vibrationally resolved ultraviolet photoelectron spectra

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Computational simulation of vibrationally resolved ultraviolet photoelectron spectra



## Structural features

- ▶ Rigid  $\sigma$ -scaffold
- ▶ two  $\pi$ -bonds (separated by the central  $\sigma$ -scaffold)

- ▶ Electronic and nuclear pieces of information are intertwined in the experimental data

- ▶ Validation of a computational approach based on a composite scheme:

electron propagator theory

+

Franck-Condon and vertical gradient approximations

- ▶ The composite scheme is expected to be effective for:
  - Semi-rigid molecular systems
  - Systems where the effects of LAMs on the experimental ultraviolet photoelectron spectrum is negligible

# Computational strategy

## Vertical gradient approximation

Duschinsky transformation:  $\overline{\mathbf{Q}} = \mathbf{J}\overline{\overline{\mathbf{Q}}} + \mathbf{K}$

- ▶ Assumptions:
  - $\mathbf{J} = \mathbf{I}$
  - $\mathbf{K} = f(\overline{\mathbf{g}})$

## Franck-Condon approximation

$$\langle \overline{x}_i | \mathcal{J}_{IF}^e | \overline{\overline{x}}_f \rangle \approx \mathcal{J}_{IF}^e(\overline{\mathbf{Q}}_{eq}) \langle \overline{x}_i | \overline{\overline{x}}_f \rangle$$

## Approximation of the electron propagator matrix

- ▶ diagonal approaches:
  - Outer Valence Green's Functions (OVGF) method
- ▶ non-diagonal approaches:
  - Non-diagonal renormalized second-order (NR2) approximation
- ▶ basis set: maug-cc-pVTZ

## Computational protocol

- ▶ 1. geometry optimization and calculation of harmonic frequencies at B3LYP/maug-cc-pVTZ level of theory
- ▶ 2. Calculation of vertical ionization energies (VIEs) at NR2/maug-cc-pVTZ level of theory
- ▶ 3. Calculation of  $\overline{\overline{\mathbf{g}}}_x$  through numerical differentiation<sup>†</sup> of VIEs calculated at OVGF/maug-cc-pVTZ level of theory
- ▶ 4. Calculation of relative intensities through a time-independent (TI) approach

<sup>†</sup>Mathematical formulations suitable for the analytical differentiation are available in literature; see:

J. V. Ortiz, *J. Chem. Phys.*, **2000**, 112, 56 - 68.

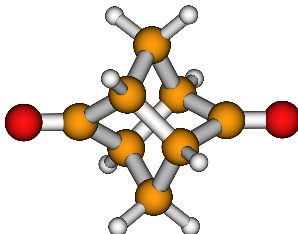
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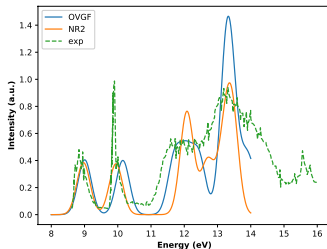
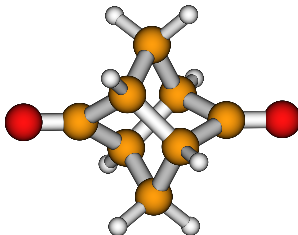
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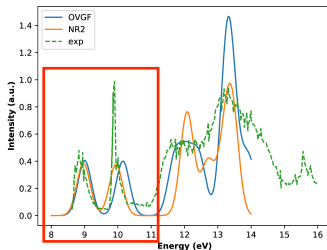
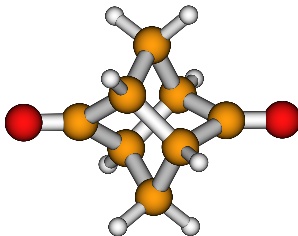
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## Comparison between calculated and experimental spectra

- ▶ Similar results with diagonal and non-diagonal approximations in the outer valence region

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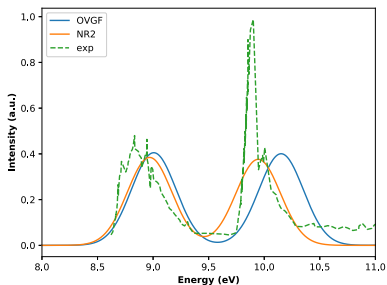
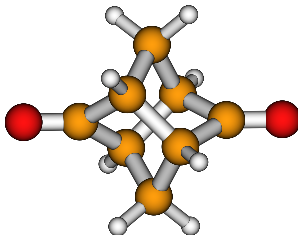
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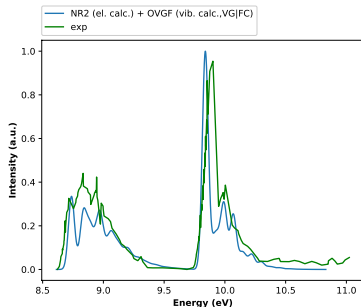
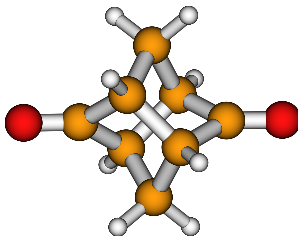
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# The case of stella-2,6-dione



## Comparison between calculated and experimental spectra

- ▶ Similar results with diagonal and non-diagonal approximations in the outer valence region
- ▶ Computation → The inclusion of vibronic of band shapes signatures is pivotal

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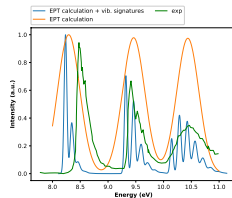
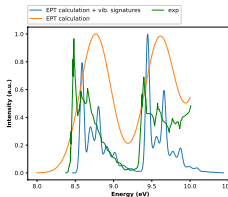
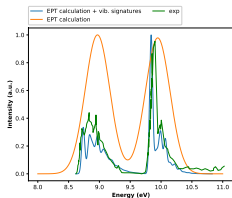
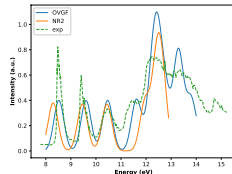
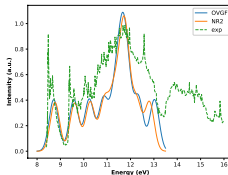
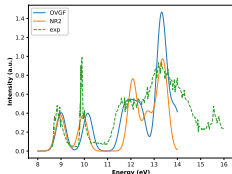
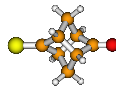
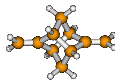
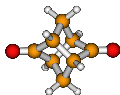
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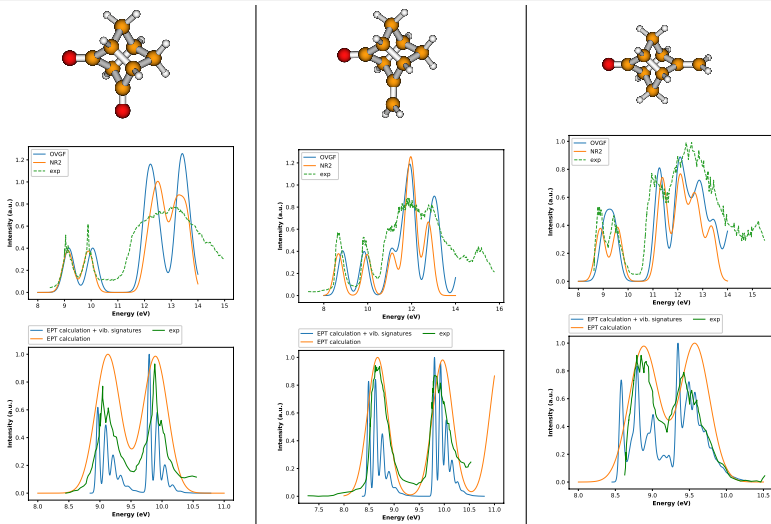
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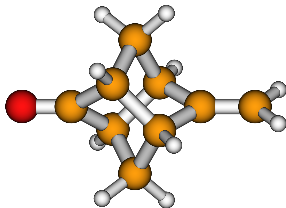
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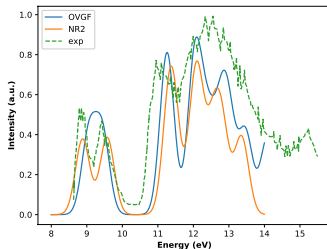
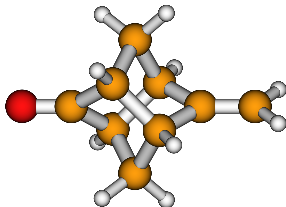
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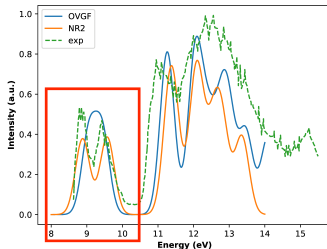
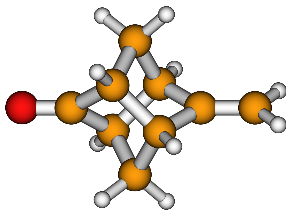
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## Comparison between calculated and experimental spectra

- Calculated VIEs at NR2 level are closer to experimental results in the outer valence region

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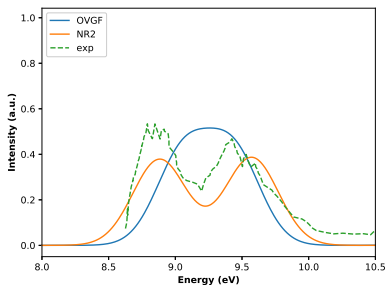
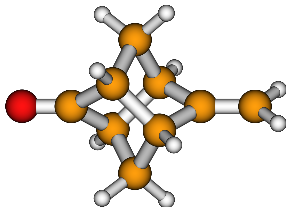
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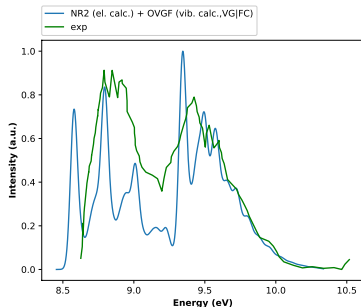
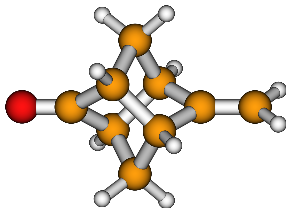
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## Comparison between calculated and experimental spectra

- ▶ Calculated VIEs at NR2 level are closer to experimental results in the outer valence region
- ▶ Disagreement between experimental and calculated bandshapes. Why?

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R. Gleiter, H. Lange, O. Borzyk, *J. Am. Chem. Soc.*, **1996**, 118, 4889 - 4895.

Computational results published in:

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# Acknowledgements

On the  
usefulness of  
electron  
propagator  
methods for a  
reliable  
computation of  
experimental  
observables

Lorenzo Paoloni

Electron  
propagator  
methods

Computational  
simulation of  
vibrationally  
resolved  
ultraviolet  
photoelectron  
spectra

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